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# Calculations of Phosphorus Nuclear Magnetic Shielding Tensors – Studies on Phosphorus Chalcogen Compounds

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The nuclear magnetic shieldings (or the NMR chemical shifts) are sensitive probes of the local electronic structure. The shielding tensors, their principal values and their principal axes system, give even more insight into the electronic structure, since they reflect the different interactions between occupied and unoccupied orbitals induced by a magnetic field oriented along the different axes.

Results of IGLO calculations<sup>[1]</sup> and experimental data for the principal values of the phosphorus shielding tensors,  $\sigma_{ii}(P)$ , within the series of  $P_4O_6Z_n$  molecules<sup>[2]</sup> (Z=O, S and Se, n=0...4) are compared in Figure 1. They agree very well, the average deviation is about 21 ppm. Even better agreement can be achieved if one determines two straight lines (with slope 1), one for  $\sigma_{33}$ , the most shielded principal value, and one for  $\sigma_{11}$  and  $\sigma_{22}$ . The intercept of the latter is some 30 ppm smaller and closer to the ideal value, 328 ppm, the shielding of the reference system. As normally found for saturated systems not having low-lying excited states (low-lying virtual orbitals) correlation effects are comparably small in these molecules. MP2-GIAO calculations<sup>[3]</sup> on the model system ZP(OH)<sub>3</sub> differ only slightly from the corresponding GIAO calculations. DFT calculations on ZP(OH)<sub>3</sub> and ZP<sub>4</sub>O<sub>6</sub> do not give any hint to large correlation effects either. They seem to be a little bit off the Hartree-Fock and the MP2 data, though.

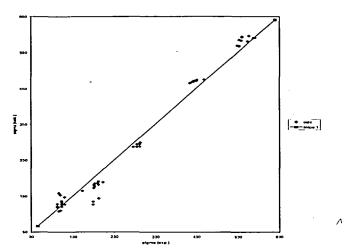


FIGURE 1 Principal values of the phosphorus shieldings,  $\sigma_{ii}(P)$ , in  $P_4O_6Z_n$ 

The principal values of  $\sigma(P)$  within the ZP<sub>4</sub>O<sub>6</sub> series are documented in Table 1. To discuss the trends observed qualitatively we recall that the shielding is given as a sum of two (or – depending on the method used – several) terms.

TABLE I MO contributions to  $\sigma_1$  and  $\sigma_2$  for the substituted phosphorus <sup>a)</sup> P4O6 OP<sub>4</sub>O<sub>6</sub> SP<sub>4</sub>O<sub>6</sub> SeP<sub>4</sub>O<sub>6</sub>  $\sigma_{\perp}$  $\sigma_{i}$  $\sigma_{i}$  $\alpha^{T}$ 189 σ 480 129 596 299 585 568 144 ΖP -385 94 -301 100 -401 19 90 -432 PO -121 -84 -113 -73 -118 -75 -120-77 <sup>a)</sup> in ppm;  $\sigma_{22} = \sigma_{11}$  because of the (local)  $C_{3v}$  symmetry; ZP – lone pair

A qualitative discussion one can be limited to the so-called paramagnetic contribution,  $\sigma^p$ . It can be expressed by the matrix elements of the first order perturbation operators and the energy difference between the occupied orbitals i and the unoccupied or virtual orbitals a. The perturbation operators are of (local) angular momentum type.

contribution for P<sub>4</sub>O<sub>6</sub>, sum of PZ bond and Z lone pairs for ZP<sub>4</sub>O<sub>6</sub>

$$\sigma_i^P \sim \langle i|\hat{l}_i|a\rangle \cdot \langle i|\hat{l}_\mu r_\mu^{-3}|a\rangle \cdot (E_a - E_i)^{-1}$$
$$l_a = (r - R_a) \times \hat{p}$$

The perpendicular component of the ZP contribution is the most deshielding one. A magnetic field perpendicular to the bond induces interactions with low-lying virtual orbitals (e.g.  $\sigma^*(PO)$ ). Since it is this bond which is formally replaced when going from one molecule to the other, it is this component that changes the most in the series. The parallel component is shielding, no deshielding interaction can be induced by a field along the bond direction. The increase of the shielding going from P<sub>4</sub>O<sub>6</sub> to the other members of the series is typically found for systems having a cylindrical charge density along an axis.

Grossmann, Potrzebowski and coworkers recently performed solid state NMR studies on <sup>71</sup>Se enriched SeP(CH<sub>3</sub>)<sub>3</sub> and SeP(C<sub>6</sub>H<sub>3</sub>)<sub>3</sub><sup>[4]</sup>. Experimental and calculated data are compared in Table 2. As found for other saturated phosphorus compounds theoretical and experimental data agree well.

TABLE 2 Principal values of σ in SeP(CH<sub>3</sub>)<sub>3</sub> and SeP(C<sub>6</sub>H<sub>5</sub>)<sub>3</sub> a)

	SeP(CH <sub>3</sub> ) <sub>3</sub>			SeP(C <sub>6</sub> H <sub>5</sub> ) <sub>3</sub> b)						
_	<b>0</b> 33	$\sigma_{22}$	$\sigma_{11}$	σ33	$\sigma_{22}$	σ <sub>11</sub>	σ33	$\sigma_{22}$	$\sigma_{11}$	
Exp.	367	287	277	376	267	224	376	267	224	
Calc.	407 381	272 286	267 285	396	239	219	397	245	216	

As pointed out above, larger deviations are observed for  $\sigma_{33}$ . Correlation effects can most probably be excluded as an explanation since MP2-GIAO calculations (for the C<sub>3v</sub> optimized structure) give only slightly different principal values than the corresponding GIAO calculations ( $\sigma_{33}$  even increases). The results of a calculation on a (SeP(CH<sub>3</sub>)<sub>3</sub>)<sub>6</sub> cluster modeling the crystal environment might be taken as an indication that intermolecular interactions are the reason for the larger deviations observed for  $\sigma_{33}$ . Further studies have to be done to check this explanation.

For low coordinate phosphorus the deshielding interactions induced by the magnetic field are distinctly larger than for the compounds discussed above. Therefore correla-

a) in ppm; see [4] for further details
b) two molecules in the unit cell, only one spectrum observed experimentally

c) data for the central molecule of a (SeP(CH3)3)6 cluster

tion effects become more important. As can be seen from Table 3 the shieldings calculated for PS<sub>3</sub><sup>-</sup> at the Hartree-Fock level converge to about -110 ppm. Inclusion of correlation effects at the MP2 level leads to a value of about 20 ppm, which is in nice agreement with recent experimental data of Karaghiosoff and Schuster<sup>[6]</sup>.

TABLE 3 Isotropic shieldings of PS<sub>3</sub><sup>-a)</sup>

INDL	7 7 1301	tobic sinc	adings of i	. 03		
IGI	LO		GIAO	Experiment		
Basis	HF	Basis	HF	MP2		
II	-93	DZP	-130	24		
III	-108	TZ2P	-114	17	31 b) 276 c)	

<sup>&</sup>lt;sup>a)</sup> in nom

PSe<sub>3</sub><sup>-</sup> is calculated to be about 70 ppm more deshielded than PSe<sub>3</sub><sup>-</sup> at the MP2 level (200 ppm at Hartree-Fock level). This is in contradiction to the experimental findings which give PSe<sub>3</sub><sup>-</sup> more shielded than PS<sub>3</sub><sup>-</sup> by 85 ppm<sup>[6]</sup>. Preliminary calculations of the spin-orbit corrections using the SOS-DFT package<sup>[7]</sup> give only a minor correction for PS<sub>3</sub><sup>-</sup> but about 160 ppm for PSe<sub>3</sub><sup>-</sup>. Correlation and spin-orbit effects have to be taken into account for PSe<sub>3</sub><sup>-</sup>, as found recently for the valence-isoelectronic CBr<sub>3</sub><sup>+ [8]</sup>.

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b) Karaghiosoff et al., ref [5]

c) Roesky et al., ref. [6]